# Size Effect of A-site Cation on n-type Thermoelectric Properties in CaMnO<sub>3</sub>-based System

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# Abstract

The thermoelectric (TE) properties of the electron-doped  $(RE_XAE_{1-X})MnO_3$  (RE: rare earth ions such as La, Pr etc., AE: the alkaline earth ions such as Ca, Sr etc., X=0.05, 0.10) have been investigated as a function of the average ionic radius  $r_A$ of the A-site cation in the perovskite structure  $(AMnO_3)$ . The electrical resistivity  $\rho$ , Seebeck coefficient S and the thermal conductivity  $\kappa$  change depending on  $r_A$ . Near the room temperature, the figure of merit Z shows a maximum at  $r_{\rm A(opt)}$ =1.182~1.184Å, which is realized for the  $La_{0.10}Ca_{0.90}MnO_3$  and  $La_{0.05}(Ca_{0.98}Sr_{0.02})_{0.95}MnO_3$ . The 5% substitution by elements M with varying valencies for the Mnsite degrades the TE properties in the  $La_XCa_{1-}$  $_X(Mn_{0.95}M_{0.05})O_3$  (X=0.05, 0.10) due to the disorder in the  $MnO_6$  network. The cracks, inevitably induced in the samples especially for the small X region, seriously reduce the TE properties.

### 1. Introduction

Since the discovery of the high thermoelectric (TE) performance in NaCo<sub>2</sub>O<sub>4</sub> [1], the TE properties of the oxides have been widely investigated. The TE module using only the p- and n-type oxides has been fabricated and operated [2]. Compared with higher p-type TE performance in NaCo<sub>2</sub>O<sub>4</sub>, the TE potential of the n-type oxides is, at present, lower for the practical use. The electron-doped  $M_XCa_{1-X}MnO_3$  shows a relatively high TE performance among oxides and is a possible candidate of the n-type TE materials. The substitution for the Ca-site by Bi has presented the best figure of merit Z for X=0.05~0.10 [3]. Ohtaki et al. investigated the 10% substitution effect for the Ca-site by several elements M in the (M<sub>0.10</sub>Ca<sub>0.90</sub>)MnO<sub>3</sub> system [4]. It was found that the electrical conductivity  $\sigma$  at 300 K increased with increasing average ionic radius of the (M<sub>0.10</sub>Ca<sub>0.90</sub>)-site resulting in M=Bi being the best element for the thermoelectricity. The origin of the high performance comes from the increase of the carrier mobility  $\mu$  caused by the larger intersite distance for hopping. However, there have been no systematic investigations as for the size effect of the A-site cation of the perovskite structure (AMnO<sub>3</sub>) on the TE properties including the Ca-site substitution by divalent ions.

In this study, we investigate the relation between the average ionic radius  $r_A$  of A-site cation and the thermoelectricity (the figure of merit Z and the power factor P) in (RE<sub>X</sub>AE<sub>1-X</sub>)MnO<sub>3</sub> (RE=rear earth ions, AE=alkaline earth ions) at the concentration of X=0.05 and 0.10. The  $r_A$  value determines the lattice distortion from the cubic through the tolerance factor  $\Gamma$  (=( $r_A$ + $r_O$ )/ $\sqrt{2}(r_{Mn}+r_O)$ ,  $r_{Mn}$  and  $r_O$  being the ionic radii of Mn and oxygen) and the one-electron bandwidth is broadened as  $\Gamma$  approaches 1. The variance  $\sigma^2$  of

the A-site ionic radius  $(=\Sigma y_i r_i^2 - r_A^2, y_i)$  being the fractional occupancy of the A-site species and  $r_i$  being its ionic radius) also affects the magnetic transition [5]. We also report the substitution effect of the Mn-site by several elements M with various ionic valencies on the thermoelectricity in La<sub>X</sub>Ca<sub>1-X</sub>(Mn<sub>0.95</sub>M<sub>0.05</sub>)O<sub>3</sub> (X=0.05 and 0.10). The relation between the thermoelectricity and the lattice distortion and the variance  $\sigma^2$  is discussed for the electron-doped system.

#### 2. Experimental

 $1.17\text{\AA} \le r_{\text{A}} \le 1.20\text{\AA}$ of The with several series RE<sub>0.10</sub>Ca<sub>0.90</sub>MnO<sub>3</sub> (RE=Dy, Gd, Sm, Nd, Pr, La), La<sub>0.10</sub>(Ca<sub>1-</sub>  $_{v}Sr_{v})_{0.90}MnO_{3}$  (0.05 $\leq$ y $\leq$ 0.4) and La<sub>0.05</sub>(Ca<sub>1-v</sub>AE<sub>v</sub>)<sub>0.95</sub>MnO\_{3} (AE=Sr, Ba:  $0.02 \le y \le 0.05$ ) were fabricated by a solid state reaction method. The  $La_{X}Ca_{1-X}(Mn_{0.95}M_{0.05})O_{3}$  (M=Mg<sup>2+</sup>, Al<sup>3+</sup>, Sn<sup>4+</sup>; X=0.05 and 0.10) samples were also fabricated. The  $r_{\rm A}$  value, the tolerance factor  $\Gamma$  and the radius variance  $\sigma^2$ were calculated from the tabulated radii with 9-fold coordination [6]. The mixtures of raw powders were calcined at 1000 °C for 24 h in air, pressed into pellets and then sintered at 1400 °C for 8 h in air. Some of the samples were heat-treated in flowing oxygen at 900 °C for 24 h in order to achieve the oxygen stoichiometry. The measured densities of samples were about 85~90 % of the ideal ones. The X-ray diffraction analyses at room temperature confirmed that all the samples were in a single orthorhombic phase. The electrical resistivity  $\rho(T)$  was measured by a standard four-terminal method from 10 K to 300 K. The thermal conductivity  $\kappa(T)$ and Seebeck coefficient S(T) were measured by a steady-state heat flow method. The carrier mobility  $\mu$  and the carrier concentration n were estimated measuring the Hall coefficient at room temperature.

#### 3. Results and Discussion

#### 3.1 (RE<sub>X</sub>AE<sub>1-X</sub>)MnO<sub>3</sub> system (X=0.05, 0.10)

Fig. 1 shows the temperature dependence of the figure of merit Z(T) of the La<sub>X</sub>Ca<sub>1-X</sub>MnO<sub>3</sub> and RE<sub>0.10</sub>Ca<sub>0.90</sub>MnO<sub>3</sub> (RE=Nd, Sm). The Z values of all the samples decrease with decreasing temperature T. In La<sub>X</sub>Ca<sub>1-X</sub>MnO<sub>3</sub>, the Z value near the room temperature takes a maximum at X=0.05 and then decreases with increasing X. The electrical resistivity  $\rho$  and the Seebeck coefficient S decrease with increasing X at room temperature, while the thermal conductivity  $\kappa$  remains nearly constant. Then the Z enhancement at X=0.05 comes from the relatively large S and small  $\rho$  values near the optimum electron doping level. It has been found that the number of cracks in the samples increases for 0≤X≤0.05 and that the existence of the cracks reduces the Z value causing the



**Figure 1:** The temperature dependence of Z(T) of the La<sub>X</sub>Ca<sub>1-X</sub>MnO<sub>3</sub> and RE<sub>0.10</sub>Ca<sub>0.90</sub>MnO<sub>3</sub> (RE=Nd, Sm).



**Figure 2:** (a)  $\rho$  and *S*, (b)  $\kappa$  and *Z* for RE<sub>0.10</sub>AE<sub>0.90</sub>MnO<sub>3</sub> at 290 K as a function of  $r_A$ .

increase in  $\rho$ . The Z value for RE<sub>0.10</sub>Ca<sub>0.90</sub>MnO<sub>3</sub> (RE=Nd, Sm) is lower than that for the RE=La sample.

Fig. 2 shows the average *A*-site ionic radius ( $r_A$ ) dependence of (a) the resistivity  $\rho$  and Seebeck coefficient *S* and (b) the thermal conductivity  $\kappa$  and figure of merit *Z* for various RE<sub>0.10</sub>AE<sub>0.90</sub>MnO<sub>3</sub> samples at 290 K. The  $\rho$  values decrease with increasing  $r_A$  for RE<sub>0.10</sub>Ca<sub>0.90</sub>MnO<sub>3</sub> and then takes a minimum at  $r_A$ =1.184Å (RE=La) and increases with further increase of  $r_A$  for the La<sub>0.10</sub>(Ca<sub>1-Z</sub>Sr<sub>Z</sub>)<sub>0.90</sub>O<sub>3</sub> system. The temperature dependence of  $\rho$  is semiconductive, *i.e.*,  $\rho$ 

increases with decreasing *T* for all the samples in this study. The *S* value hardly depends on  $r_A$  and the  $\kappa$  value slightly decreases with increasing  $r_A$ . The increase of  $r_A$  results in the reduction of the static lattice distortion from the cubic structure with the tolerance factor  $\Gamma$  approaching 1 as shown in the upper axis of abscissas. The  $\rho$  reduction in RE<sub>0.10</sub>Ca<sub>0.90</sub>MnO<sub>3</sub> may result from the decrease in the lattice distortion from the cubic, originating from the broadening of the bandwidth. The  $\rho$  increase for  $r_A$ >1.184 Å may come from the introduction of the lattice disorder and/or the increase of the radius variance  $\sigma^2$  caused by the substitution of Sr for the Ca-site.

The similar  $r_A$  dependences can be seen for the 5% electron doping system. Fig. 3 shows (a) the resistivity  $\rho$ , (b) the Seebeck coefficient S and (c) the thermal conductivity  $\kappa$ for  $(RE_{0.05}AE_{0.95})MnO_3$  samples at 290 K as a function of  $r_A$ . In these figures, the bold symbols show the values of which the samples were carefully cut from the nearly crack-free region of the sintered pellets. For the majority samples bearing some cracks as shown in the open symbols, the  $\rho$ value takes a minimum and the S value takes a faint maximum around  $r_{A} \sim 1.183$  Å for La<sub>0.05</sub>(Ca<sub>0.98</sub>Sr<sub>0.02</sub>)<sub>0.95</sub>MnO<sub>3</sub>. The  $\kappa$ value slightly decreases with increasing  $r_A$ . These behaviors are similar to those for the RE<sub>0.10</sub>AE<sub>0.90</sub>MnO<sub>3</sub> shown in Fig. 2. For the nearly crack-free samples, the low  $\rho$  value and the high  $\kappa$  value are naturally expected. The extremely strong enhancement in the S value was observed for the La<sub>0.05</sub>(Ca<sub>0.98</sub>Sr<sub>0.02</sub>)<sub>0.95</sub>O<sub>3</sub> and La<sub>0.05</sub>Ca<sub>0.95</sub>MnO<sub>3</sub> samples. One possible explanation is that the intrinsic higher TE potential was masked out by the existence of the cracks.



**Figure 3:** The  $r_A$  dependence of (a)  $\rho$ , (b) S and (c)  $\kappa$  for RE<sub>0.05</sub>AE<sub>0.95</sub>MnO<sub>3</sub> at 290 K.



**Figure 4:** The  $r_A$  dependence of Z for RE<sub>0.05</sub>AE<sub>0.95</sub>MnO<sub>3</sub> at 290 K.



**Figure 5:** The *Z* value for the RE<sub>0.05</sub>AE<sub>0.95</sub>MnO<sub>3</sub> samples at 290 K as a function of  $\sigma^2$ .

Fig. 4 shows the figure of merit Z as a function of  $r_A$  for RE<sub>0.05</sub>AE<sub>0.95</sub>MnO<sub>3</sub> samples at 290 K. For both the crackbearing and nearly crack-free samples, the Z value takes a maximum around  $r_A$ =1.812~1.814 Å, depending on the  $r_A$ value. The Z value obtained in this study was lower than the reported values of Bi<sub>0.04</sub>Ca<sub>0.96</sub>MnO<sub>3</sub> (Z=98 x 10<sup>-6</sup> K<sup>-1</sup>) [7] and Bi<sub>0.10</sub>Ca<sub>0.90</sub>MnO<sub>3</sub> (Z=350 x 10<sup>-6</sup> K<sup>-1</sup>) [4].

Fig. 5 shows the Z value for RE<sub>0.05</sub>AE<sub>0.95</sub>MnO<sub>3</sub> at 290 K as a function of the A-site ionic radius distribution (variance)  $\sigma^2$ . The Z value takes a maximum and then decreases with increasing  $\sigma^2$ . The increase of  $\sigma^2$  indicates the increase of the A-site lattice randomness in the perovskite structure. In Fig. 4, the Zincrease with increasing  $r_A$  for  $r_A < 1.184$  Å may result from the decrease in the lattice distortion and the Z decrease for  $r_A > 1.184$  Å may result from the increase of  $\sigma^2$  and/or the introduction of the crystal disorder caused by the substitution by Sr or Ba ions for the Ca-site.

Fig. 6 shows the mobility  $\mu$  and the carrier concentration *n* for the RE<sub>0.05</sub>AE<sub>0.95</sub>MnO<sub>3</sub> samples at 290 K as a function of  $r_A$ , which were estimated by the Hall coefficient. The mobility  $\mu$  shows a maximum at  $r_A \sim 1.184$  Å, where the  $\rho$ 



**Figure 6:** The mobility  $\mu$  and the carrier concentration *n* for RE<sub>0.05</sub>AE<sub>0.95</sub>MnO<sub>3</sub> at 290 K as a function of  $r_A$ .

shows a lower value shown in Fig. 3(a). The *n* value, at which *Z* takes a maximum, is  $1 \times 10^{21}$  cm<sup>-3</sup> and these results indicate that the carrier hopping is the main conduction mechanism in this system. However, it must be noted that only the average radius  $r_A$  is not the main factor to determine  $\mu$  and  $\sigma$  (=1/ $\rho$ ) as clarified the maximum at  $r_A \sim 1.184$  Å.

## 3.2 La<sub>X</sub>Ca<sub>1-X</sub>(Mn<sub>0.95</sub>M<sub>0.05</sub>)O<sub>3</sub> system (X=0.05, 0.10)

We have investigated the effect of the Co-site substitution by the various elements on the TE properties for La<sub>0.80</sub>Sr<sub>0.20</sub>(Co<sub>1-Z</sub>M<sub>Z</sub>)O<sub>3</sub> [8,9]. It was found that the ionic valency of the substituted element M more strongly influenced the thermoelectricity rather than the ionic radius or spin state of the element M. The valence effect of the dopants for the Mn-site substitution on the TE properties were studied for  $La_XCa_{1-X}(Mn_{0.95}M_{0.05})O_3$  (X=0.05, 0.10; M=Mg<sup>2+</sup>, Al<sup>3+</sup>, Sn<sup>4+</sup>). The substitution by the 2+ ions (Mg<sup>2+</sup>) is expected to make the average Mn valency increase and the substitution by the 4+ ions (Sn<sup>4+</sup>) is expected to make the average Mn valency slightly decrease. Fig. 7 shows (a) the resistivity  $\rho$ , (b) the Seebeck coefficient S, (c) the thermal conductivity  $\kappa$  and (d) the figure of merit Z for  $La_XCa_{1-X}(Mn_{0.95}M_{0.05})O_3$  (X=0.05, 0.10) at 290 K. The  $\rho$  values increase and the  $\kappa$  values decrease for the Mn-site substitution owing to the introduction of the disorder in the MnO<sub>6</sub> network. The  $\rho$  values for the Al<sup>3+</sup> substitution are the largest and are about an order of magnitude higher than those of the non-substituted samples. The absolute values of S slightly increase by the Mn-site substitution for X=0.10, but decrease for X=0.05. As a result, the Z value decreases by the Mn-site substitution for both X=0.05 and 0.10 samples relative to the non-substituted La<sub>x</sub>Ca<sub>1-x</sub>MnO<sub>3</sub>. The Mn-site substitution seriously damages the TE properties irrespective of the ionic valency of the substituted elements in contrast with the La<sub>0.80</sub>Sr<sub>0.20</sub>(Co<sub>1-</sub>  $_{Z}M_{Z}O_{3}$  system.

#### 4. Conclusions

The thermoelectric (TE) properties of the electron-doped  $(RE_XAE_{1-X})MnO_3$  (X=0.05, 0.10) have been investigated as a



**Figure 7:** The X dependence of (a)  $\rho$ , (b) *S*, (c)  $\kappa$  and (d) *Z* of La<sub>X</sub>Ca<sub>1-X</sub>(Mn<sub>0.95</sub>M<sub>0.05</sub>)O<sub>3</sub> at 290 K (M=Mg<sup>2+</sup>, Al<sup>3+</sup>, Sn<sup>4+</sup>).

function of the average ionic radius  $r_A$  of the A-site cation in the perovskite structure. The electrical resistivity  $\rho$ , Seebeck coefficient S and the thermal conductivity  $\kappa$  have been measured and the figure of merit Z and the power factor P have been estimated. Important experimental results and conclusions obtained in this study are summarized as follows.

- (1) The  $\rho$ , *S* and  $\kappa$  values depend on  $r_A$ . Near the room temperature, the *Z* value shows a maximum at  $r_{A(opt)}=1.182\sim1.184$ Å, which is realized for the La<sub>0.10</sub>Ca<sub>0.90</sub>MnO<sub>3</sub> and La<sub>0.05</sub>(Ca<sub>0.98</sub>Sr<sub>0.02</sub>)<sub>0.95</sub>MnO<sub>3</sub>. The decrease of *Z* with decreasing  $r_A$  from 1.182 Å may result from the decrease of the tolerance factor  $\Gamma$  and the decrease of *Z* with increasing  $r_A$  from 1.184 Å may result from the introduction of the crystal disorder and/or the increase of the radius variance  $\sigma^2$ . The *P* value shows the similar  $r_A$  dependence.
- (2) The 5% substitution by elements M for the Mn-site degrades the TE properties in  $La_XCa_{1-X}(Mn_{0.95}M_{0.05})O_3$  (X=0.05, 0.10) compared with the non-substituted samples. The TE property degradation of this system is mainly caused by the crystal disorder in the MnO<sub>6</sub> network and is nearly independent of the ionic valency of the substituted elements.
- (3) The cracks, inevitably induced in the samples during the sintering process, seriously reduce the TE properties especially for the small X region (0≤X≤0.10).

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