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## Trapped magnetic field properties of MgB<sub>2</sub> bulks doped with Ti

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### Abstract

We have studied the trapped magnetic field properties on doped MgB<sub>2</sub> bulk with Ti (nominal content was Mg:Ti =0.9:0.1 and 0.8:0.2). The trapped magnetic field,  $B_T$ , was enhanced by approximately 1.3 times from 2.6 T of the pristine bulk to 3.5 T of the Ti-doped bulks at 15 K. The extrapolated  $B_T(T)$  curve reached 5 T below 4.2 K, indicating that 5 T class bulk magnet can be realized using MgB<sub>2</sub> with Ti-doping. The critical current density,  $J_c$ , under the magnetic field was also enhanced by Ti-doping, therefore, an irreversibility field,  $B_{irr}$ , of the Ti-doped bulks exceeded 5 T at 20 K. Ti and TiB<sub>2</sub> impurities confirmed by the powder X-ray diffraction acted as the pinning centers and resulted in the enhancement of the trapped field and the critical current density.

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**Keywords:** MgB<sub>2</sub>; superconducting bulk magnet; Ti doping

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### 1. Introduction

A large polycrystalline MgB<sub>2</sub> superconducting bulk with centimeter size in diameter and in thickness has been fabricated to realize a quasi-permanent magnet because of weak-link-free grain boundaries (GBs) (Kambara et al. (2001)). However, since the bulk fabricated by an *in-situ* method under an ambient pressure has intrinsically a low packing density of about 50% (Naito et al. (2012)), the high-pressure sintering is one of the routes to enhance the density and the critical current density,  $J_c$ , i.e., the trapped magnetic field,  $B_T$ . Just after the discovery of MgB<sub>2</sub>, MgB<sub>2</sub> bulk sintered under a pressure of 2 GPa offered the possibility of 2 T class bulk magnets at 6 K (Viznichenko et al. (2003)). We also reported the 2.5 T MgB<sub>2</sub> bulk magnet operated at 12.7 K using dense bulk, which was prepared by a hot isostatic pressing (HIP) method under a pressure of 98 MPa (Sasaki et al. (2013)). Since the GBs are well known to act as pinning centers in MgB<sub>2</sub>, the effect of the grain-size refining on the flux pinning has been investigated. Recently, Fuchs *et al.* achieved a trapped field of 5.4 T in hot pressed MgB<sub>2</sub> using ball-milled powder, which demonstrates that simultaneous grain refining and densification are extremely effective to enhance  $B_T$  (Fuchs et al. (2013)). Contrary to the mechanical milling, the impurity doping such as Ti was also reported to fine down the MgB<sub>2</sub> grains. Zhao *et al.* argued that the critical current density was improved over three orders of magnitude from  $\sim 10^3$  A/cm<sup>2</sup> of the pristine sample to  $\sim 10^6$  A/cm<sup>2</sup> of Ti10% doped sample at 20 K (Zhao et al. (2001)). However,

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further Ti-doping suppressed the  $J_c$  value and the  $J_c$  of the Ti20% sample was about one order of magnitude smaller than that of the Ti10% sample at 20 K.  $\text{TiB}_2$  with lamellar structure was created between nanometer-size grains and acted as the flux pinning center, which was observed by TEM (Zhao et al. (2002)). Therefore, fine grains, high packing density and a novel flux pinning center such as  $\text{TiB}_2$  are also possible to bring about a higher  $J_c$  and higher  $B_T$ . In this paper, we report on the fabrication of  $\text{MgB}_2$  bulks doped with two Ti contents by the HIP method in order to realize the high density and the introduction of novel pinning center at the same time. We measured their trapped field and critical current density, and finally discuss the effect of Ti-doping.

## 2. Experimental

$\text{MgB}_2$  bulk samples were fabricated by an *in-situ* HIP method. The precursors were prepared as follows; raw powders of Mg (99% in purity), Ti (99% in purity), and amorphous B (99% in purity) were weighted with  $(1-x):x:2.0$  in molar ratio and ground for 1 h. The mixture was pressed into a pellet 40 mm in diameter and 20 mm in thickness under an uniaxial pressure of about 12 MPa. Subsequently the pellet was further densified by a cold isostatic pressing method under a pressure of 196 MPa. The precursor pellet was sealed in a stainless steel container under vacuum by an electron beam welding method. The  $\text{MgB}_2$  precursor in the container was sintered at 900 °C for 3 h under Ar-gas pressure of 98 MPa in the HIP furnace and cooled down to room temperature by furnace cooling. The disc-shaped  $\text{MgB}_2$  bulk with about 38 mm diameter and 6.9 mm thickness was prepared by the dry machining.

The  $\text{MgB}_2$  bulk was magnetized by field cooling (FC) in a magnetic field of 5 T parallel to the thickness direction using a 10 T cryogen-free superconducting magnet (JMTD-10T100, Japan Superconductor Technology (JASTEC), Inc.) and then the applied magnetic field was decreased to 0 T at a rate of 0.022-0.22 T/min. The bulk was placed on the cold-stage of a helium refrigerator by inserting a thin indium sheet to obtain good thermal contact. The trapped field was measured by a cryogenic Hall sensor (BHT-921, F.W. Bell Inc.) mounted on the center of the bulk surface. The temperature of the bulk was monitored by a Cernox thermometer which was adhered beside the Hall sensor. After the FC magnetization experiments, we evaluated the critical current density  $J_c$ , using a small piece cut from the bulk.  $J_c$  was estimated from the magnetic hysteresis using the extended Bean model (Bean (1962)) following  $J_c = 20\Delta M/a(1 - a/3b)$ , where  $\Delta M$  is the width of the hysteresis, and  $a$  and  $b$  ( $a < b$ ) are the cross-sectional dimensions of the sample perpendicular to the applied magnetic field. The magnetic hysteresis was measured using a commercial SQUID magnetometer (MPMS-XL, Quantum Design Inc.).

## 3. Results and discussion

Fig. 1 shows the temperature dependence of the trapped magnetic field  $B_T(T)$  of the pristine and 10% and 20% Ti doped bulk pieces. The  $B_T$  value of the pristine bulk is about 2.46 T at the lowest temperature of 17 K, decreases with increasing temperature, and finally reaches zero at around 39 K, which originates from the decrease of  $J_c$  with increasing  $T$ . The highest  $B_T$  of 3.45 T at 15 K was achieved for the Ti10% bulk; which was approximately 1.3 times larger than the extrapolated value of 2.6 T for the pristine bulk at 15 K. The Ti20% bulk offers almost the same  $B_T$  profile as that of the Ti10% bulk. The  $B_T(T)$  curve extrapolated to lower temperature reaches 5 T around 4.2 K for Ti10% bulk, which allowed us to realize a 5 T class superconducting bulk magnet. It is noteworthy that the critical temperature  $T_c$ , defined by the temperature at  $B_T=0$ , is independent of the Ti content. The results suggest that Ti is not a substitute for the Mg-site, because the site substitution usually deteriorates the superconducting properties.

Fig. 2 shows the magnetic field dependence of the critical current density  $J_c(B)$  of the pristine and Ti-doped bulks at 20 K. The  $J_c(B)$  curves of both Ti-doped bulks are nearly the same, which is consistent with the  $B_T(T)$  behavior. The  $J_c$  value of the pristine bulk decreases monotonically with increasing magnetic field, especially a strong decay is observed under the applied field higher than 4 T. Although the  $J_c$  in self-field is not so enhanced by Ti-doping, but in magnetic field it is rather improved. For instance, the  $J_c=4.1\times 10^3$  A/cm<sup>2</sup> of the pristine sample is increased to  $2.0\times 10^4$  A/cm<sup>2</sup> under the applied field of 3 T by 10% Ti-doping. An irreversibility field  $B_{ir}$ , defined by the magnetic field at  $J_c=10$  A/cm<sup>2</sup>, obviously exceeds 5 T in the Ti10% sample. Assuming the uniform  $J_c$  distribution in the bulk, the average macroscopic  $J_c$ , which creates the trapped field, can be estimated by the following analytical equation

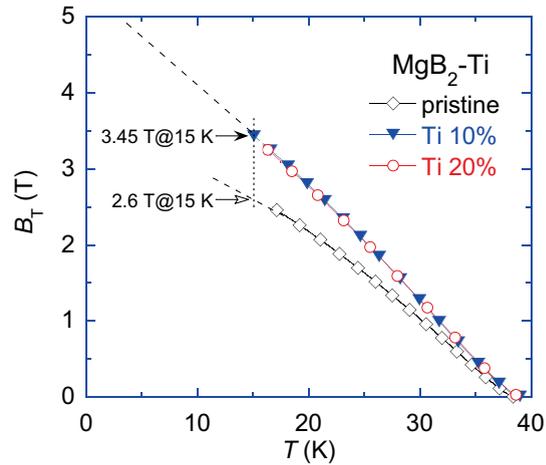


Fig. 1: Temperature dependence of the trapped magnetic field of the pristine and Ti-doped MgB<sub>2</sub> bulks. The dotted lines represent the extrapolation lines.

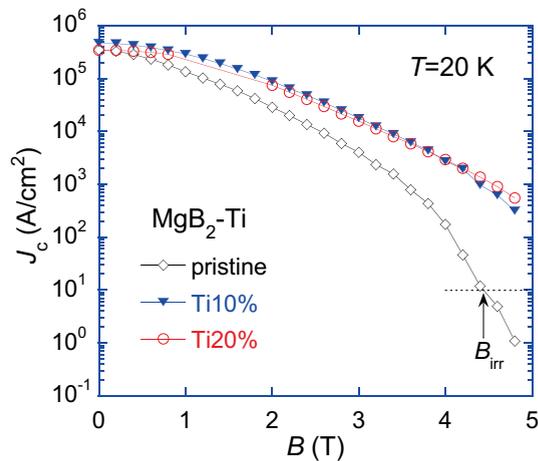


Fig. 2: Magnetic field dependence of the critical current density at 20 K for the pristine and Ti-doped MgB<sub>2</sub> bulks. An irreversibility field,  $B_{irr}$ , is defined by a criterion of  $J_c=10$  A/cm<sup>2</sup>.

based on the Bio-Savart law,

$$B_T(z) = \frac{\mu_0 J_c}{2} \left( (z+t) \ln \frac{(D/2) + \sqrt{(D/2)^2 + (z+t)^2}}{z+t} - z \ln \frac{(D/2) + \sqrt{(D/2)^2 + z^2}}{z} \right), \quad (1)$$

where,  $D$  and  $t$  are the diameter and thickness of the disc and  $z$  is the distance from the center of the top surface of the bulk. The trapped field of 2.17 T for the pristine bulk is enhanced to 2.79 T by the Ti10% doping at 20 K, as shown in Fig. 1. Using Eq. (1) and the measured  $B_T$  values, the average macroscopic  $J_c$  is estimated to be  $2.9 \times 10^4$  A/cm<sup>2</sup> for the pristine and  $3.7 \times 10^4$  A/cm<sup>2</sup> for the Ti10% bulk, which correspond to the measured  $J_c$  under the applied field of 2.0 T and 2.6 T, respectively. These results suggest that the experimentally obtained  $B_T$  values can be roughly reproduced by the Bio-Savart law.

Fig. 3 shows the powder X-ray diffraction (XRD) pattern taken at room temperature for the pristine and Ti-doped bulks. The XRD peaks of the MgB<sub>2</sub> phase are labeled by “MgB<sub>2</sub>”, and an impurity phase is not found for the pristine

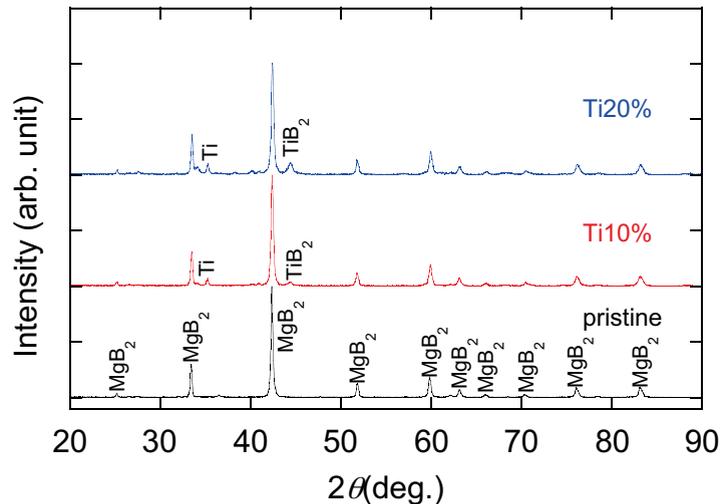


Fig. 3: Powder X-ray diffraction patterns taken at room temperature for the pristine and Ti-doped  $\text{MgB}_2$  bulks.

bulk. On the other hand, the impurity phases such as Ti and  $\text{TiB}_2$  appear for the Ti-doped bulks, especially the peak intensity due to  $\text{TiB}_2$  increases obviously with the increase in the Ti contents. Ti contained impurities act as the pinning center, and enhance both the  $J_c$  and  $B_T$ . If the grain of  $\text{MgB}_2$  is fined down to nanometer size by Ti-doping, as pointed out by Zhao *et al.* (Zhao et al. (2001, 2002)), the trapped field is expected to be over 5 T at around 12 K as observed in the hot-pressed bulk using ball-milled powder Fuchs et al. (2013). However, the highest  $B_T$  of 3.45 T at 15 K in the Ti-doped bulks does not reach the expected value. Although the Ti-doping successfully certainly introduced the flux pinning centers, the refining effect of the grain size was not be confirmed at the present stage. Microscopic observation to clarify the grain size is now in progress.

#### 4. Summary

We have studied the effects of Ti-doping on the trapped magnetic field and critical current density of the  $\text{MgB}_2$  bulks. Two bulk samples with the nominal composition of Mg:Ti=0.9:0.1 and 0.8:0.2 have been measured. The trapped magnetic field was considerably increased by Ti-doping; the highest value of 3.45 T achieved for both Ti-doped bulks at 15 K was about 1.3 times larger than that of the pristine bulk. The trapped field line extrapolated to lower temperature reached 5 T below 4.2 K. The critical current density was also enhanced by Ti-doping and the irreversibility field, which was defined by the criterion of  $J_c=10 \text{ A/cm}^2$ , exceeded 5 T at 20 K. Ti and  $\text{TiB}_2$  phases confirmed by the X-ray diffraction acted as the flux pinning centers. It is noteworthy to point out that both trapped field and critical current density obtained for the Ti-doped bulk pieces hardly depended on the Ti contents. Further studies such as microstructure analysis and further Ti-doping must be considered.

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