

## Trapped field and temperature rise in MgB<sub>2</sub> bulks magnetized by pulsed field

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The pulsed field magnetization (PFM) was performed for the MgB<sub>2</sub> bulk 30 mm in diameter by a capsule method, and the applied pulsed-field dependence of the trapped field  $B_z$ , temperature rise  $\Delta T$  and time dependence of the local field  $B_L^C(t)$  were measured on the bulk.  $B_z=0.71$  T was achieved at 16 K after applying the magnetic pulsed field of  $B_{ex}=1.55$  T, where  $\Delta T_{max}$  of about 10 K took place. The results were compared with those by field-cooled magnetization (FCM), and the magnetic flux intrusion and flux trap for the MgB<sub>2</sub> bulk during PFM are discussed.

### INTRODUCTION

The superconducting bulk magnet using a REBaCuO (RE: rare earth element or Y) is one of the exemplary models for practical applications such as a sputtering cathode, magnetic separation and drag delivery system [1], which produces tesla-order quasi-permanent magnets. However, the large single-domain bulk over 100 mm in diameter was difficult to fabricate. MgB<sub>2</sub> bulk has also a promising potential as quasi-permanent magnet, which has several attractive natures for bulk superconducting magnet, such as low-cost and light-weight, which are clear contrast with the REBaCuO bulk magnet. The problem of weak-links at the grain boundaries can be ignored in the MgB<sub>2</sub> polycrystalline bulk due to their long coherence length,  $\xi$  [2]. The characters enable us to realize the better and larger polycrystalline MgB<sub>2</sub> bulk magnets below the superconducting transition temperature  $T_c=39$  K. Several groups have already reported the trapped field on the MgB<sub>2</sub> bulk by the field-cooled magnetization (FCM), and attained the trapped field over 1.5 T at low temperatures [3-5].

A pulsed-field magnetization (PFM) has also been investigated to magnetize the bulk superconductors because of an inexpensive and mobile experimental set-up with no need for a superconducting magnet. However, for the REBaCuO bulks, the trapped field  $B_z$  achievable by PFM is nonetheless lower than that achievable by FCM because of the large temperature rise caused by the dynamical motion of the magnetic flux. Several approaches have been performed and succeeded in enhancing  $B_z$  using the multi-pulse techniques [6, 7]. We have experimentally examined the trapped field  $B_z$  on the surface of cryocooled REBaCuO bulks during PFM for various starting temperatures  $T_s$  and applied fields  $B_{ex}$  [8]. To enhance  $B_z$ , the reduction in temperature rise and the lowering of  $T_s$  are effective. Considering the obtained experimental results, we proposed a new PFM technique named a modified multi-pulse technique with stepwise cooling (MMPSC) and successfully realized a highest field trap of  $B_z=5.20$  T on a GdBaCuO bulk with 45 mm in diameter at 30 K [9], which is a record-high value by PFM to date.

In this paper, we applied the PFM technique to the large MgB<sub>2</sub> bulk fabricated by a capsule method [10]. The PFM experiments for the MgB<sub>2</sub> bulk were performed for the first time and the results were compared with those by FCM.

### EXPERIMENTAL PROCEDURE

MgB<sub>2</sub> bulks were fabricated by the *in-situ* capsule method. The detailed procedure of the sample preparation was described elsewhere [10]. Raw powders of Mg (99% in purity) and amorphous B (99% in purity) were weighted with 1.1 : 2.0 in molar ratio, ground and was pressed into pellet 30 mm in diameter

and 9 mm in thickness under uni-axial pressure in air. The precursor pellet sealed in the capsule under Ar gas atmosphere was sintered at 800 °C for 6 h in a box furnace and cooled to room temperature.

The MgB<sub>2</sub> bulk situated in the stainless steel flange of the capsule was tightly anchored onto the cold stage of a Gifford–McMahon (GM) cycle helium refrigerator. The initial temperature  $T_s$  of the bulk was 14 K. The magnetizing solenoid coil (94 mm I.D., 153 mm O.D., and 67 mm height), which was dipped in liquid nitrogen, was placed outside the vacuum chamber. A magnetic pulse  $B_{ex}$  up to 1.85 T with a rise time of 0.01 s and a duration of 0.15 s was applied to the bulk by flowing the pulsed current from a condenser bank. The time evolutions of the local field  $B_L^C(t)$  and the subsequent trapped field  $B_z$  at the center of the bulk surface were monitored by the Hall sensor (BHT 921; F W Bell) using a digital oscilloscope, which was adhered the center of the bulk surface. Two-dimensional trapped field profiles of  $B_z(1\text{ mm})$  were mapped on the bulk with 1 mm above the bulk surface, stepwise with a pitch of 1 mm by scanning the axial-type Hall sensor (BHA 921; F W Bell) using an  $x$ - $y$  stage controller. During PFM, the time dependence of temperature  $T(t)$  was also measured at the bulk surface of using the Cernox™ (Lakeshore Cryotronics, Inc.) thermometer.

FCM was also performed for the MgB<sub>2</sub> bulk using a cryo-cooled superconducting solenoid magnet (JMTC-10T100, JASTEC, Inc.). Under the magnetic field of 5 T, the bulk was cooled to  $T_s$  between 16 K to and 30 K, and then the applied field was reduced to zero with a speed of  $-3\text{ mTs}^{-1}$ . The trapped field  $B_T^C$  at the center of the bulk surface was measured by the Hall sensor (BHT 921; F W Bell) adhered at the center of the bulk surface. The trapped field profiles,  $B_T^{FCM}$  (4 mm) were measured at  $T_s=28\text{ K}$  on the vacuum sheath surface at a distance of  $z=4\text{ mm}$  from the bulk surface.

## RESULTS AND DISCUSSION

Figure 1(a) shows the trapped field  $B_z=B_T^C$  at the center of the bulk surface, as a function of the strength of the applied pulsed field  $B_{ex}$ .  $B_T^C$  increases for  $B_{ex}>1\text{ T}$ , takes a maximum at  $B_{ex}=1.54\text{ T}$  and then decreases with increasing  $B_{ex}$ . The maximum  $B_T^C$  was 0.71 T. The  $B_T^C$  vs  $B_{ex}$  curve is a typical one for PFM. Figure 1(b) shows the time dependences of the temperature change on the bulk surface for each applied field  $B_{ex}$ . The temperature increases abruptly just after the pulse application and sharply decreases with increasing time. The magnitude of the maximum temperature rise increases with increasing  $B_{ex}$ . The sharp temperature change results from the low specific heat and high thermal conductivity of the MgB<sub>2</sub> bulk, compared with those for the REBaCuO bulk at operating temperatures.

Figure 2 presents the trapped field profiles  $B_z$  on the MgB<sub>2</sub> bulk 1 mm above the bulk surface for various applied pulsed fields  $B_{ex}$ . For  $B_{ex}=1.04\text{ T}$  as shown in Fig. 2(a), the magnetic flux was trapped at the right lower region of the bulk, where the critical current density  $J_c$  seems to be relatively lower in the region. For  $B_{ex}=1.23\text{ T}$  as shown in Fig. 2(b), the magnetic flux was also trapped at the left upper region of the bulk. For  $B_{ex}=1.42\text{ T}$  and  $1.54\text{ T}$  as shown in Figs. 2(c) and 2(d), the trapped field profiles are nearly the conical one. The trapped field profiles look like nearly homogeneous, compared with those for the REBaCuO bulks [8]. Because the MgB<sub>2</sub> bulk was fabricated by the *in-situ* sintering method. On the other hand, the REBaCuO bulks was fabricated by the melt-growth, in which the four-fold growth sector boundaries (GSBs) existed and the inhomogeneous  $J_c$  distribution was shown in the bulk.

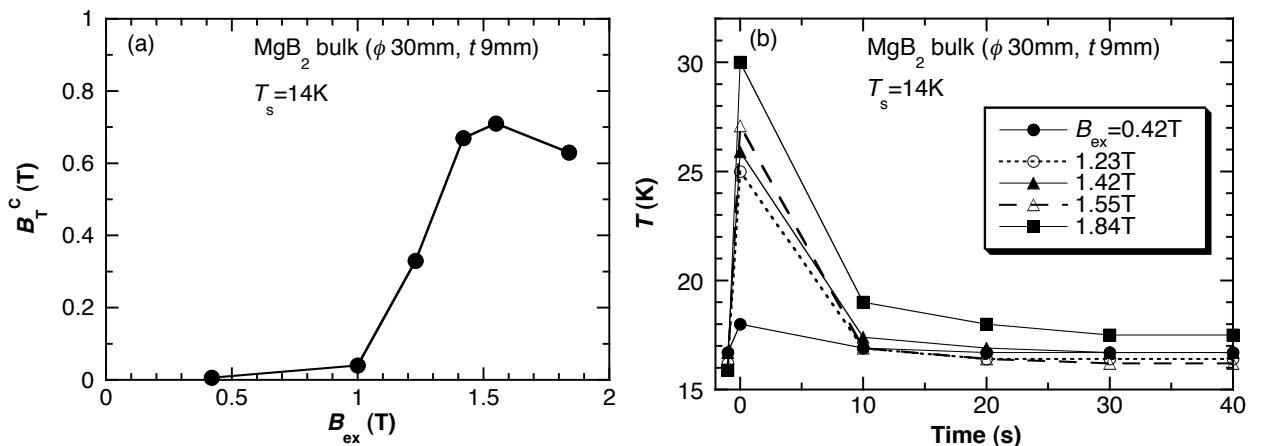


Figure 1 (a) The trapped field  $B_T^C$  at the center of the bulk surface as a function of applied pulsed field  $B_{ex}$  on the MgB<sub>2</sub> bulk at  $T_s=14\text{ K}$ . (b) The time dependences of the temperature change on the bulk surface for each applied field  $B_{ex}$ .

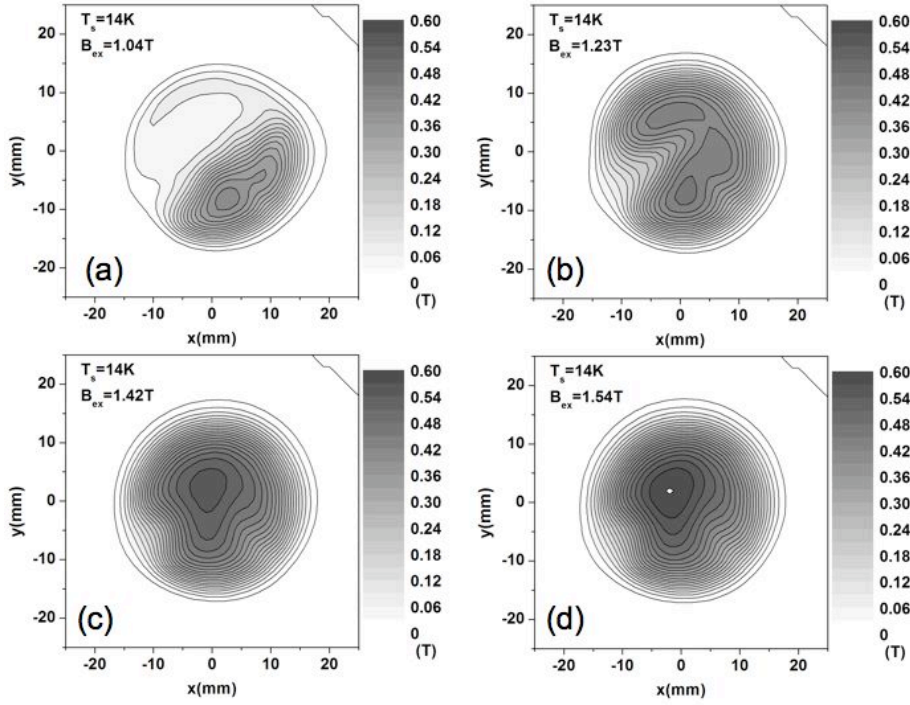


Figure 2 The trapped field profiles  $B_z$  on the MgB<sub>2</sub> bulk 1 mm above the bulk surface at (a)  $B_{ex}=1.04$  T, (b) 1.23 T, (c) 1.42 T and (d) 1.54 T.

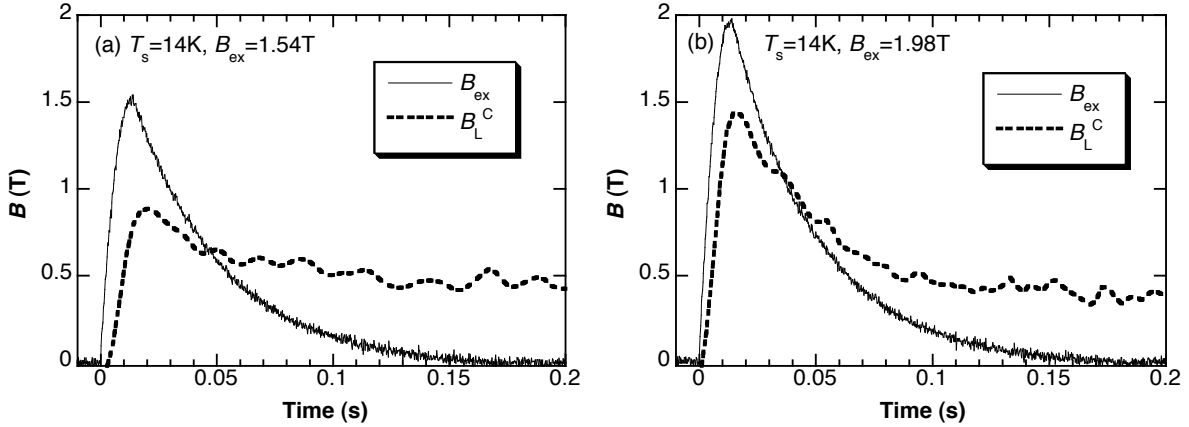


Figure 3 Time dependences of the local field  $B_L^C(t)$  and the applied pulsed field  $B_{ex}(t)$  for (a)  $B_{ex}=1.54$  T and (b) 1.98 T. The waving in  $B_L^C$  comes from the smoothing process.

Figure 3 shows the time dependences of the local field  $B_L^C(t)$  after applying the pulsed field of  $B_{ex}=1.54$  T and 1.98 T. Time dependence of pulsed field  $B_{ex}(t)$  is also shown. For each  $B_{ex}$ ,  $B_L^C(t)$  starts to increase with a slight time delay, takes a maximum at 0.015 s and then decreases to a final value due to the flux flow. The maximum  $B_L^C(t)$  increases with increasing  $B_{ex}$  and is smaller than that of  $B_{ex}(t)$  because of the shielding effect in the superconductor.

Figure 4(a) depicts the trapped field  $B_T^C$  by FCM at the center of the bulk surface as a function of temperature.  $B_T^C$  increases with decreasing temperature, which originates from the enhancement of  $J_c$  at low temperatures. The highest  $B_T^C$  value was 1.50 T at 16 K, which was smaller than that for the reported MgB<sub>2</sub> bulk with 28 mm in diameter fabricated under high pressure [5]. This might come from the fact that the density (approximately 1.3–1.4 g/cm<sup>3</sup>) of the present bulk is lower than that of the bulk in Ref. 5.

Figure 4(b) presents the trapped field profile  $B_T^{FCM}$  (4 mm) on the bulk magnetized by FCM at  $T_s=28$  K. The maximum  $B_T^{FCM}$  (4 mm) was 0.32 T. The trapped field profile is conical, which suggests that the  $J_c$  distribution in the bulk is nearly homogeneous from the view point of the profile by FCM. For practical applications, the magnetic field over 0.3 T can be used on the vacuum chamber at present, which is lower than that of the Nd-Fe-B permanent magnet. The MgB<sub>2</sub> bulk magnet over 3 T on the bulk surface and over 1 T on the vacuum sheath can be realized by the enhancements of  $J_c$ , the mass density and the connectivity between the grains in the polycrystalline MgB<sub>2</sub> bulk.

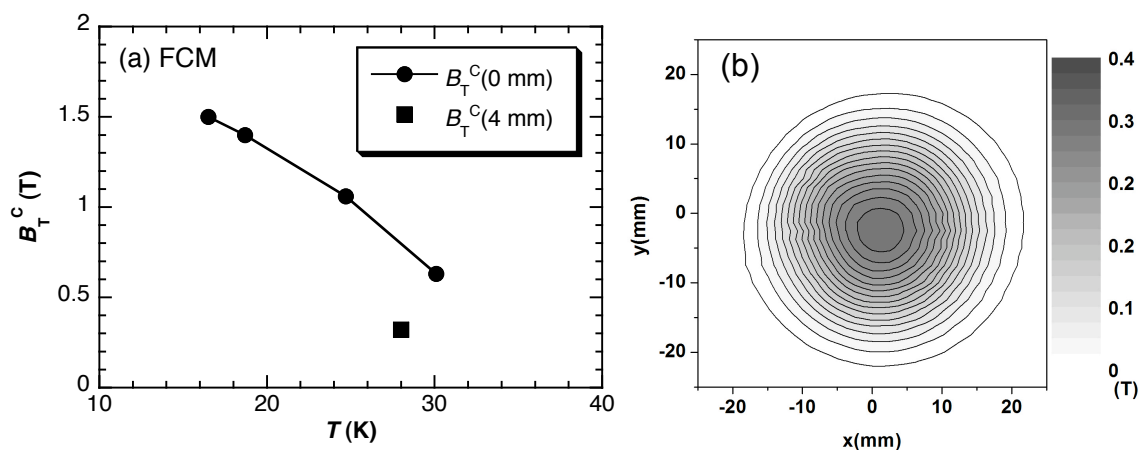


Figure 4 (a) The trapped field  $B_T^C$  by FCM at the center of the bulk surface as a function of temperature  $T_s$ . (b) The trapped field profile  $B_T^{\text{FCM}}$  (4 mm) on the bulk magnetized by FCM at  $T_s=28$  K. The maximum  $B_T^{\text{FCM}}$  (4 mm) was 0.32 T.

## SUMMARY

The pulsed-field magnetization (PFM) and the field-cooled magnetization (FCM) were performed for the MgB<sub>2</sub> bulk 30 mm in diameter fabricated by a capsule method, and the applied pulsed field dependence of the trapped field  $B_z$ , temperature rise  $\Delta T$  and time dependence of the local field  $B_L^C(t)$  were measured during PFM on the MgB<sub>2</sub> bulk.  $B_z=0.71$  T was trapped by PFM at 14 K after applying the magnetic pulse of  $B_{\text{ex}}=1.55$  T, where  $\Delta T_{\text{max}}$  of about 10 K took place.  $B_T=1.5$  T was trapped by FCM at 16 K at the center of the bulk surface with a conical trapped field distribution. PFM is a promising technique to magnetize the MgB<sub>2</sub> bulk besides the FCM.

## ACKNOWLEDGMENT

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