Effect of Metal Ring Setting Outside HTSC Bulk Disk on Trapped Field and Temperature Rise in Pulse Field Magnetizing

Hiroyuki Fujishiro, Kazuya Yokoyama, Masahiko Kaneyama, Tetsuo Oka, and Koshichi Noto

Abstract—In order to enhance the trapped field in cryo-cooled HTSC bulk magnets using pulse field magnetizing (PFM), a metal ring (stainless steel 304 and/or Al) has been tightly set onto the SmBaCuO bulk disk and the relation between the total trapped flux \( \Phi_T \), the trapped field \( B_T \), and the temperature rise \( \Delta T \), has been investigated as a function of the applied pulse field \( B_{ex} \). The \( \Phi_T \) and \( B_T \) values are enhanced about 10\% - 20\% by the metal ring due to the reduction in the temperature rise \( \Delta T \). These results suggest that a part of the generated heat \( Q \) due to the flux motion in the peripheral region promptly transfers to the metal ring and the heat transfer to the cold stage is improved by the ring setting.

Index Terms—High \( T_c \) bulk superconductors, metal ring setting, pinning and viscous loss, pulse field magnetizing, temperature measurement.

I. INTRODUCTION

In view of the practical applications of high-\( T_c \) bulk superconductors (HTSCs) as a high strength bulk magnet for a magnetic levitation system and so on, pulse field magnetizing (PFM) as well as the static field-cooled magnetizing (FCM) has been intensively investigated because of the relatively compact and inexpensive setup. The field trapped by PFM is, however, lower than that attained by FCM at temperatures below 77 K. The main reason has been attributed to the large heat generation due to the dynamic motion of the magnetic flux against the vortex pinning force \( F_p \) and the viscous force \( F_v \). We have studied the temperature rise \( \Delta T \) and trapped field \( B_T \) on the surface of cryo-cooled YBaCuO and SmBaCuO bulk disks during PFM and pointed out the importance of the \( \Delta T \) reduction for the \( B_T \) enhancement [1]–[4]. The total generated heat \( Q \) was estimated using the specific heat \( C \) of the bulk and the maximum \( \Delta T \). It was pointed out that the \( B_T \) values as functions of the initial temperature \( T_0 \) and applied pulse field \( B_{ex} \) can be understood on the basis of the trapped field \( B_T \) by FCM vs. the temperature \( T \) diagram, i.e., by the decrease of the critical current density \( J_c \) associated with the temperature rise. In order to enhance the \( B_T \) and \( \Phi_T \) values by PFM, the reduction of the \( \Delta T \) is an indispensable issue. The iteratively magnetizing pulsed-field method with reducing amplitude (IMRA) [5], locating yoke pieces around a bulk [6], a use of vortex-type coils [7] and a multi-pulse technique with step-wise cooling (MPSC) [8] have been attempted to suppress the heat generation during PFM. A \( B_T \) of 17.24 T at 29 K in YBaCuO bulk has been realized by FCM, enhancing the thermal conduction of the bulk by impregnating a high thermal conductivity alloy into drilled holes [9].

In this study, we set a metal ring (Al and/or stainless steel 304) tightly on a SmBaCuO bulk disk. The following effects are expected from the metal ring set on to the HTSC bulk disk. First, since the heat generation due to PFM occurs mainly in the peripheral region, the metal ring attached to the bulk disk periphery can easily and promptly receive the heat. Second, since the \( ab \)-plane thermal conductivity \( \kappa_{ab} \) is far larger than the \( c \)-axis \( \kappa_c \), the generated heat can easily reach the metal ring along the radial paths of the disk, then transferring to the cold stage through the highly conductive ring. These two effects should contribute to the \( \Delta T \) reduction of the bulk. Finally, the metal ring setting helps to mechanically reinforce the bulk disk [10].

II. EXPERIMENTAL PROCEDURE

A highly \( c \)-axis oriented SmBaCuO bulk superconductor (Dowa Mining Co., Ltd) with 45 mm diameter and 15 mm thickness was used. This consisted of four growth sector regions (GSRs), divided by the growth sector boundaries (GSBs) [2]. The bulk is composed of \( \text{SmBa}_2\text{Cu}_3\text{O}_y \) (Sm123) and \( \text{SmBa}_2\text{Ba}_3\text{Cu}_3\text{O}_{y} \) (Sm211) with the molar ratio of Sm123 : Sm211 = 1.0 : 0.3, 15.0 wt.% A\( g_2 \)O powder, and 0.5 wt.% Pt powder. The bulk was uniformly impregnated by epoxy resin in vacuum and then the epoxy resin layer on the surface of the bulk disk was removed. The stainless steel (SUS304) or the aluminum (Al) ring with 4 mm in thickness and 15 mm in height, was fixed onto the bulk disk using apiezon-N grease. We call the ring-attached sample as SUS-Sm or Al-Sm. Since a large heat generation due to the eddy current took place in the Al ring \( \Delta T > 6 \text{ K at 100 K} \), the ring was cut with a gap of 1 mm as shown in Fig. 1. Resultantly, \( \Delta T \) was diminished to within 0.2 K. Since the Al ring did not contact tightly with the bulk due to the gap, another SUS ring...
with 1 mm in thickness was set outside the Al-ring. Hereafter, we denote this sample with the Al+SUS (W) ring the W-Sm. Table I summarizes the sizes, the volume of each metal ring, and the specific heat, the heat capacity (J/K) at 40 K of the Sm-bulk and each metal ring, and the ratio of the heat capacity of the metal ring to that of the Sm-bulk. The heat capacity increases by 30%, 18% and 27% for the SUS-ring, Al-ring and W-ring, respectively, relative to that for the bare Sm bulk. The \( \kappa \)-plane and \( c \)-axis thermal conductivity \( \kappa \) at 40 K of the Sm-bulk and that of SUS and Al are also shown [11]. \( C_v \) and \( \kappa \) were measured for the present specimens [12].

The bulk disk was tightly stacked on the sapphire plate (45 mm in diameter and 20 mm in thickness) attached to the cold stage of a helium refrigerator. The initial stage temperature was kept at 40 K. The temperatures, \( T_0 \) at the center of the bulk (P0) and, \( T_1 \sim T_4 \) at P1 \sim P4 were monitored by fine chromel-constantan thermocouples adhered to the upper bulk surface by GE7031 varnish. P1 \sim P4 were situated on the central radial lines of each GSR by 9 mm apart from P0. The temperature of the metal ring \( TR \) was also measured at PR. The total trapped magnetic flux \( \Phi_T \) and the distribution of the trapped magnetic flux density \( B_T^{3\text{mm}} \) were measured using an axial type Hall sensor (F.W. Bell, model BHA 921), which scanned 3 mm above the bulk surface stepwise with a pitch of 1.2 mm. The trapped field \( B_T^{3\text{mm}} \) on the bulk surface was measured by the Hall sensor adhered to the position at PH with a 2.5 mm distance from P0. Five iterative magnetic pulses (No1 \sim No5) with the same amplitude \( B_{ex} \) from 3.83 T to 6.04 T (rise time: 12 ms) were applied sequentially after re-cooling to \( T_s \). \( T(t) \), \( \Phi_T^{3\text{mm}} \) and \( B_T^{3\text{mm}} \) were measured at each stage.

### III. RESULTS AND DISCUSSION

Fig. 2(a) shows the time dependence of temperature \( T_3(t) \) at P3 after applying the No1 and No5 pulse of \( B_{ex} = 4.70 \text{ T} \). For the bare Sm-bulk (B-Sm) without the ring, the magnetic

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**TABLE I**

<table>
<thead>
<tr>
<th></th>
<th>Sm-bulk</th>
<th>SUS-ring</th>
<th>Al-ring</th>
<th>W-ring (Al-SUS)</th>
</tr>
</thead>
<tbody>
<tr>
<td>I.D. (mm φ)</td>
<td></td>
<td>45</td>
<td>45</td>
<td>45-53</td>
</tr>
<tr>
<td>O.D. (mm φ)</td>
<td></td>
<td>45</td>
<td>45</td>
<td>53-55</td>
</tr>
<tr>
<td>Height (mm)</td>
<td>15</td>
<td>15</td>
<td>15</td>
<td>15-15</td>
</tr>
<tr>
<td>( V ) (mm³)</td>
<td>23.84</td>
<td>9.24</td>
<td>9.24</td>
<td>9.24-2.54</td>
</tr>
<tr>
<td>( C_v ) (J/cm³K) at 40 K</td>
<td>0.48</td>
<td>0.38</td>
<td>0.23</td>
<td>-</td>
</tr>
<tr>
<td>( C ) (J/K) at 40 K</td>
<td>11.44</td>
<td>3.47</td>
<td>2.08</td>
<td>3.04</td>
</tr>
<tr>
<td>( C/\kappa ) ratio</td>
<td>1.0</td>
<td>0.30</td>
<td>0.18</td>
<td>0.27</td>
</tr>
<tr>
<td>( \kappa ) at 40K (mW/cmK)</td>
<td>130 (ab)</td>
<td>55</td>
<td>-5000</td>
<td>-</td>
</tr>
<tr>
<td></td>
<td>30 (c)</td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

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Fig. 1. Photograph of the SmBaCuO-bulk disk fitted with the W(Al+SUS)-ring. The positions of the temperature and magnetic field measurements (P0 \sim P4, PR and PH) are shown.
flux preferentially intrudes into the bulk from the GSRs containing P2, P3 and P4 and the sharp and large temperature rises ($\Delta T = 25$ K) take place within 3 s at these positions. $T(t)$ recovers to the initial temperature $T_s$ after $15\sim 20$ min. $T_3(t)$ of SUS-Sm decreases faster after the $T_3(t)$ maximum, although the maximum temperature rise is similar to that of B-Sm. This result means that a part of the generated heat $Q$ promptly transfers to the SUS-ring. The maximum $\Delta T$ value at $T_3$ decreases to $\sim 20$ K for Al-Sm and W-Sm. The rate of temperature decrease ($-d\Delta T(t)/dt$) after the peak is small for Al-Sm but that for the W-Sm is large and almost the same as the SUS-Sm. As for the iterative pulse field applications, $\Delta T$ is the largest for the No1 pulse and decreases for the succeeding pulses [2]. The increment of the trapped field $\Delta B_T$ is also the largest for the No1 pulse, followed by a gradual increase for the No2 and No3 pulses. For the No5 pulse application, the $T(t)$ peak disappears for all the settings. $\Delta T$ is slightly smaller for the SUS-Sm and W-Sm than for the bare bulk and Al-ring setting.

Fig. 2(b) presents the $T(t)$ at $T_0$ for each metal ring setting. $T(t)$ rises up latest among $T_0 \sim T_4$ because the distance from the major heat source is usually the longest [2]. For the SUS-Sm and W-Sm, $T(t)$ reaches a maximum faster than that for B-Sm. These results mean that the flux motion is enhanced in the central region of the bulk by the SUS-Sm and W-Sm and the magnetic fluxes are accumulated faster in the bulk center. The anomalously long time constant for the temperature rise for the Al-ring setting may result from the high thermal contact resistance between the bulk and Al-ring.

Fig. 2(c) shows the $T_R(t)$ for each metal ring setting. The maximum $\Delta T$ values for SUS-Sm and W-Sm are 12 K and 10 K, respectively, and the time constant for the temperature rise of $T_R(t)$ is longer for the SUS-Sm and W-Sm than that of $T_3(t)$. The temperature rise in the metal ring is mainly due to the heat conduction from the bulk and the temperature rise in the Al-ring is slower because of the high thermal contact resistance.

Figs. 3(a) and (b) summarize the maximum temperature rise of the bulk $\Delta T_{B_{\text{max}}}$ and the ring $\Delta T_{R_{\text{max}}}$ after the No1 and No5 pulse as a function of $B_{\text{ex}}$. $\Delta T_{B_{\text{max}}}$, $\Delta T_{R_{\text{max}}}$, the averaged value at $\Delta T_{B_{\text{max}}} \sim \Delta T_{R_{\text{max}}}$, increases with increasing $B_{\text{ex}}$ for both No1 and No5 pulses. $\Delta T_{B_{\text{max}}}$ of B-Sm is the largest and is reduced by the metal ring setting. Especially, the reduction of the $\Delta T_{B_{\text{max}}}$ is the largest for W-Sm and Al-Sm; e.g., from 23 K to 18 K for the No1 pulse of $B_{\text{ex}} = 4.70$ T. $\Delta T_{R_{\text{max}}}$ also increases with increasing $B_{\text{ex}}$ for both No1 and No5 pulses and is about 50 to 80% of $\Delta T_{B_{\text{max}}}$. $\Delta T_{R_{\text{max}}}$ is the largest for the SUS-ring and the smallest for the W-ring.

Figs. 4(a) and 4(b) present the total trapped flux $\Phi_{T}\text{max}$ for each metal ring setting after the No1 and No5 pulse as a function of $B_{\text{ex}}$. For the No1 pulse, $\Phi_{T}\text{max}$ of B-Sm is smallest for each $B_{\text{ex}}$ and the $\Phi_{T}\text{max}$ value is enhanced by the metal ring setting. Especially, for the SUS-Sm and W-Sm, the $\Phi_{T}\text{max}$ values after the 4.70 T and 5.53 T pulses increase about 20% compared with those of the bulk without ring. The increase may mainly come from the decrease of the temperature rise due to the ring setting, which results in the increase of the effective critical current density. The $\Phi_{T}\text{max}$ enhancement is observed for Al-Sm, which suggests that the good thermal contact between the ring and the Sm-bulk is of vital importance. The trapped field $B_{T}\text{max}$ measured at PH is also enhanced by the metal ring setting; for example, $B_{T}\text{max} = 2.95$ T for W-Sm is about 10% larger than that of B-Sm ($= 2.71$ T) after the No5 pulse of $B_{\text{ex}} = 5.53$ T. The typical trapped field distributions $B_{T}\text{max}$ are shown in the inset of Fig. 4(a) in which a conical field distribution can be confirmed.

Neglecting the heat drained to the cold stage, the generated heat $Q$ by the PFM operation can be estimated using the following equation,

\[
Q = \int_{T_s}^{T_s+\Delta T_{B_{\text{max}}}} (C_{\text{Bulk}} V_{\text{Bulk}}) dT + \int_{T_s}^{T_s+\Delta T_{R_{\text{max}}}} (C_{\text{Ring}} + V_{\text{Ring}}) dT,
\]

where $C_{\text{Bulk}}$ and $C_{\text{Ring}}$ are the specific heat (J/cm$^3$K), and $V_{\text{Bulk}}$ and $V_{\text{Ring}}$ are the volume of the Sm-bulk and the metal ring, respectively. Fig. 5 shows the estimated $Q$ values as a function of $B_{\text{ex}}$ for each metal ring setting. For the No1 pulse, the $Q$ value of the bulk with each metal ring is smaller than that of B-Sm. This result suggests that the heat transferred from the bulk is partially drained to the cold stage through the ring. For
bulk. The good thermal conductance between the bulk periphery and ring is necessary to reduce $\Delta T$ and to enhance the $\Phi_T^P$ and $B_T^P$ values.

(3) The metal ring with a high thermal conductivity also enhances the thermal drain for the generated heat $Q$ to transfer to the cold stage.

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REFERENCES


